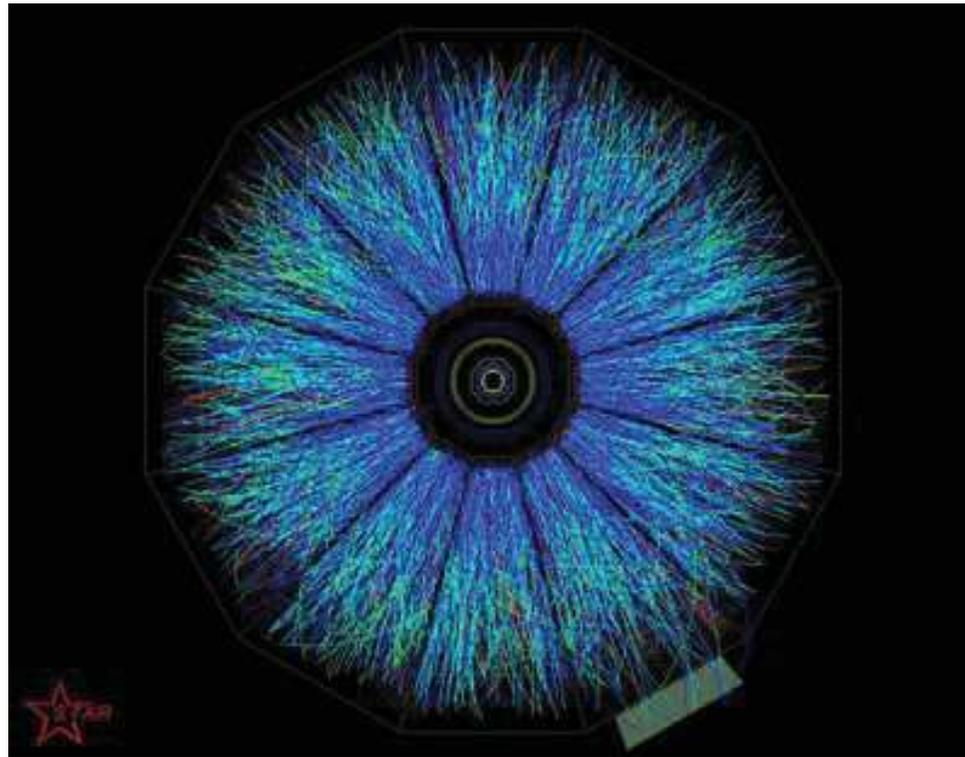


# 7. Luminosity

## ▷ Luminosity

Gaussian beams. Hourglass effect. Space charge limit.  
Luminosity optimization. Tune spread. Beam-beam limit.  
Beam separation. Crab crossing. Crab sextupoles. Beam disruption

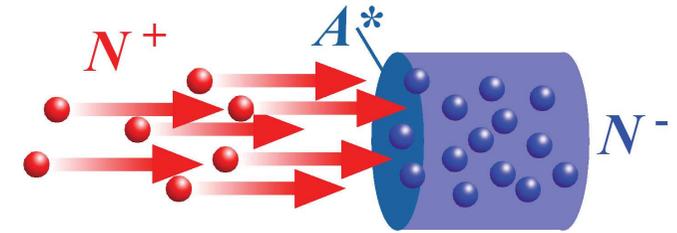


# Luminosity with Gaussian beams

Luminosity = particles/time  $\times$  particle/area

$$\mathcal{L} = \frac{N^+}{T} \times \frac{N^-}{A^*}$$

( $A^*$  common interaction area)



Luminosity = 4-d overlap of particle distributions  $\varrho^\pm$  ( $\beta = 1$ ):

$$\mathcal{L} = f_c \int \int \int \int_{-\infty}^{+\infty} \varrho^+(x, y, s + ct) \varrho^-(x, y, s - ct) 2c dt ds dx dy$$

$2c$  = relative velocity of bunches in laboratory system,

$f_c = c/b$  = collision frequency,  $b$  distance between successive bunches.

Gaussian distributions  $\varrho^\pm$ , also include [horizontal] crossing angle  $2\theta \ll 1$ :

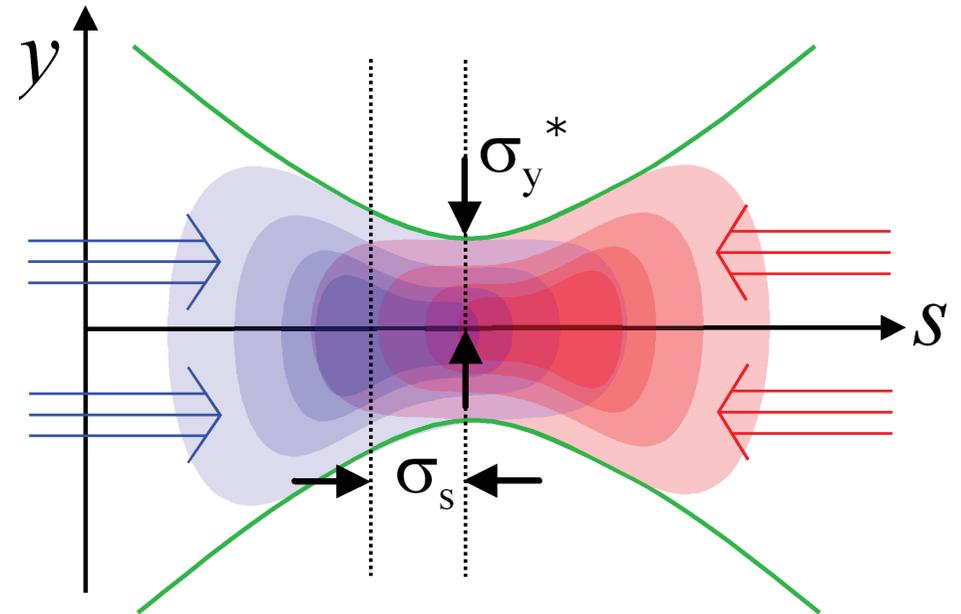
$$\varrho^\pm(x, y, s \pm ct) = \frac{N^\pm}{(2\pi)^{3/2} \sigma_x(s) \sigma_y(s) \sigma_s} e^{-\frac{(x \pm s\theta)^2}{2\sigma_x(s)^2} - \frac{y^2}{2\sigma_y(s)^2} - \frac{(s \pm ct)^2}{2\sigma_s^2}}.$$

Focus at collision point (\*):  $\sigma_u(s) = \sigma_u^* \sqrt{1 + \left(\frac{s}{\beta_u^*}\right)^2}$ ,  $u = x; y$

$$\int \Rightarrow \mathcal{L} = \frac{f_c N^+ N^-}{4\pi \sigma_x^* \sigma_y^*} \cdot S \quad \longrightarrow \quad A^* = 4\pi \sigma_x^* \sigma_y^* \quad \text{for Gaussian beams}$$

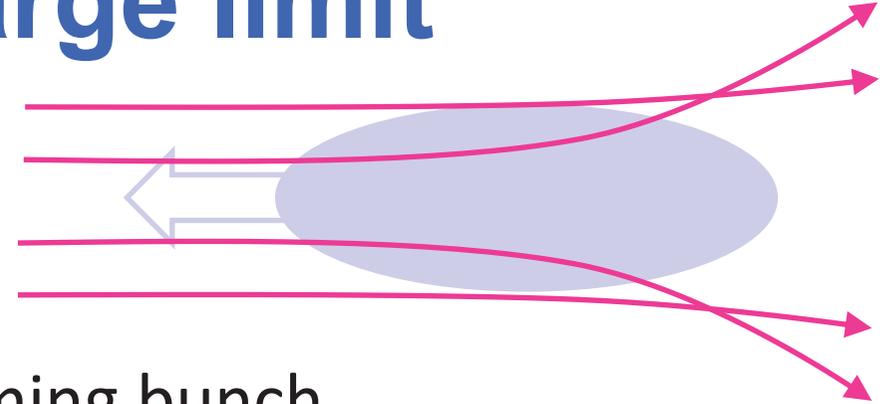
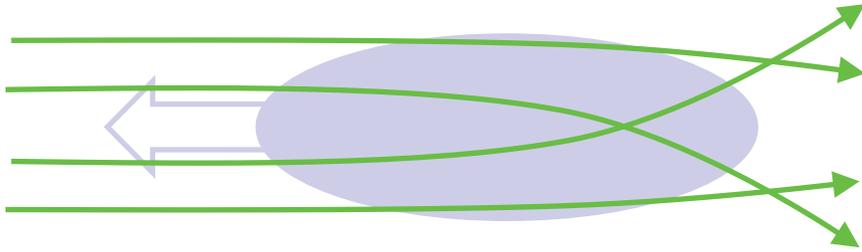
Luminosity suppression factor  
*hourglass effect*

$$S = \frac{2}{\sqrt{\pi} \sigma_s} \int_0^\infty \frac{e^{-\left(\frac{s}{\sigma_s}\right)^2} e^{-\left(\frac{\theta s}{\sigma_x(s)}\right)^2}}{\sqrt{1 + \left(\frac{s}{\beta_x^*}\right)^2} \sqrt{1 + \left(\frac{s}{\beta_y^*}\right)^2}} ds$$



Limit on focus:  $\beta^* > \sigma_s \quad \longrightarrow \quad S \approx 0.8 \dots 0.95$

# Space charge limit



“Lens” formed by charge in oncoming bunch

- ◆ very strong: focal length  $f \sim \text{cm} !$
- ◆ focusing (e.g.  $e^+e^-$ ) or defocusing (e.g.  $pp$ )
- ◆ non-linear
  - strong for core particles ( $x, y$  small)
  - weak for halo particles ( $x, y$  large)
  - even stronger for some  $|\Delta s| > 0$  particles

Storage ring collider: betatron tunes  $Q_x, Q_y$

→ Beam collision: tune shifts  $\Delta Q_{x,y}(x, y, \Delta s) \Rightarrow$  tune spread

→ **Beam-beam tune shift parameter** or **space charge parameter**:  
tune shift of center particle  $\zeta_{x,y} = \Delta Q_{x,y}(0,0,0)$

# Beam-beam limit

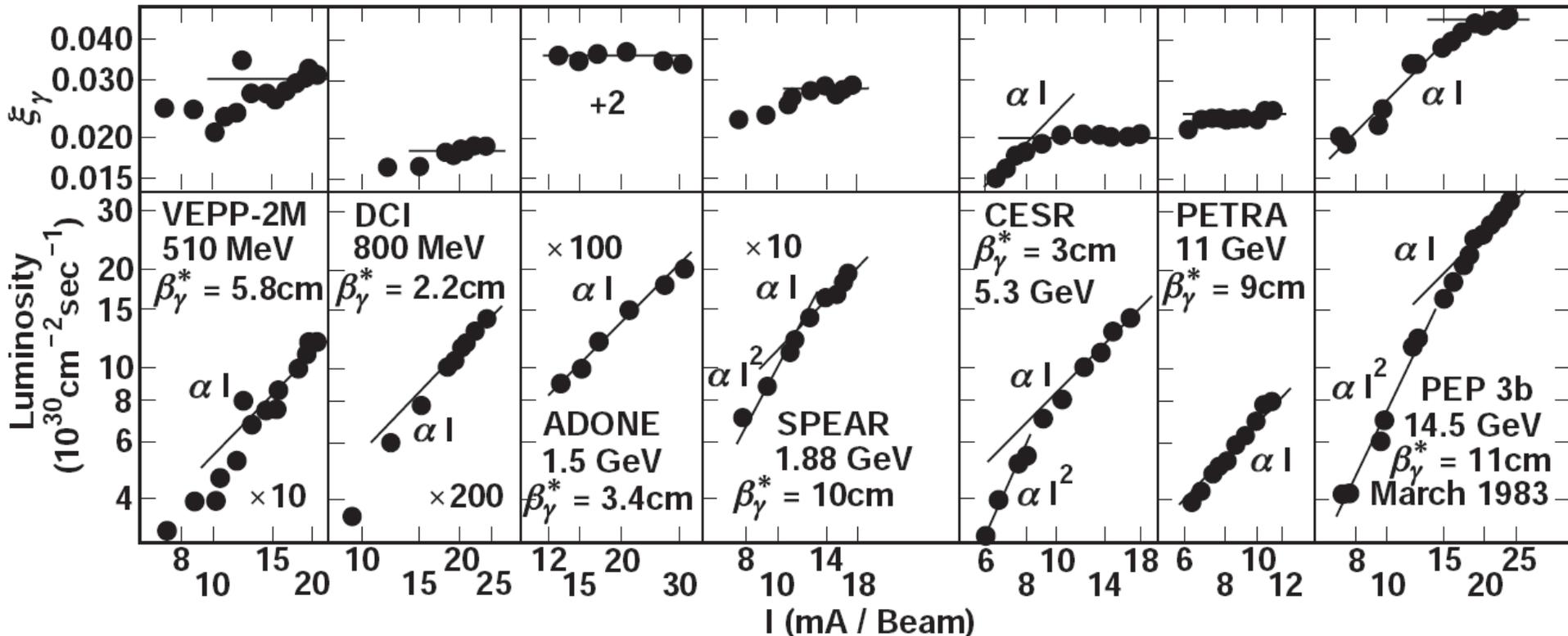
Empirical limits for space charge parameter

$$\zeta < 0.05 \text{ in } e^+e^- \quad \zeta < 0.005 \text{ in } p\bar{p} \text{ collisions}$$

→ Saturation of  $\zeta$  with current

→ blow-up of beams, increase of interaction area

→ only linear increase of luminosity:  $\mathcal{L} \sim N$ , not  $\mathcal{L} \sim N^2$



R. Talman, Specific luminosity limits of  $e^+e^-$  colliders, PRSTAB 5, 081881 (2002)

# Tune spread

Non-linearity of beam-beam lens  $\implies$  Tune spread (*beam's footprint*)

Force vs. position:

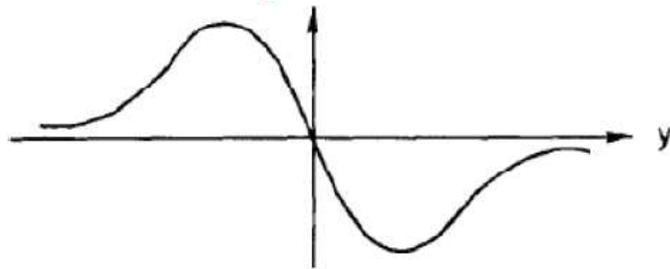
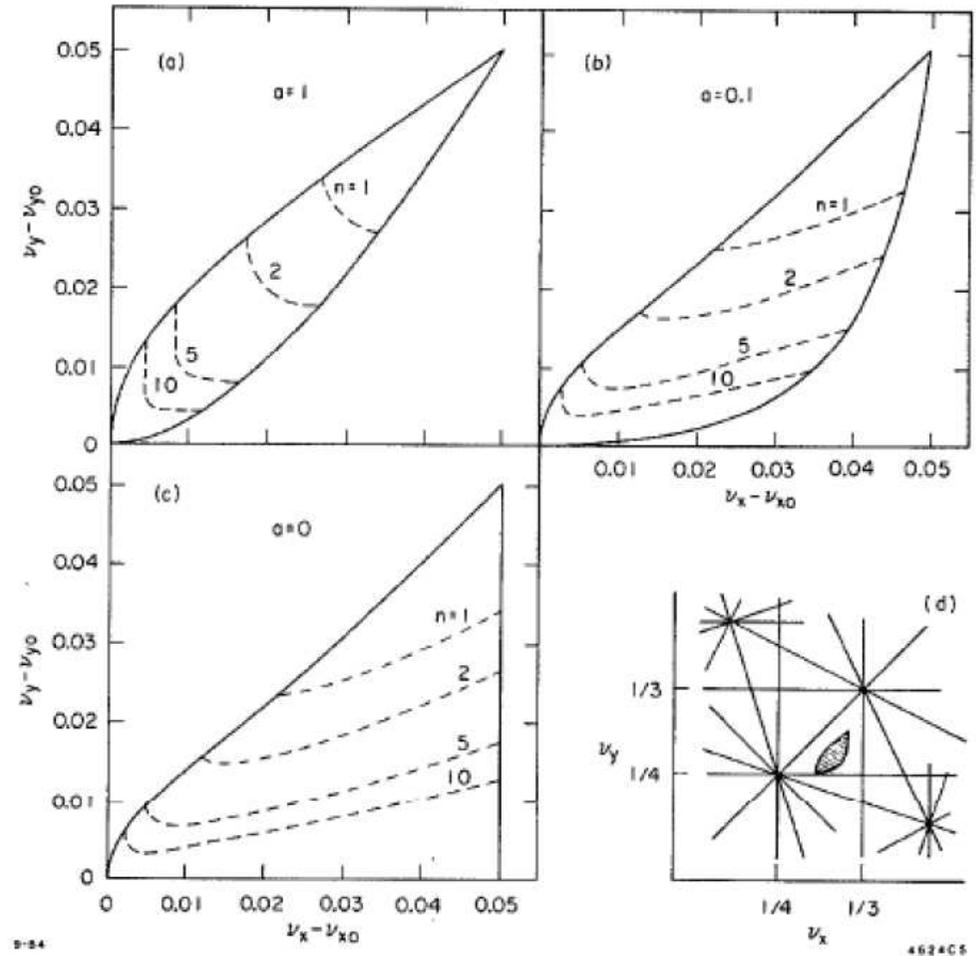


Fig. 11. Beam-beam tune spreads. We assume the two beams have opposite charges.  $(\nu_{x0}, \nu_{y0})$  is the unperturbed working point. With beam-beam collisions, the working point extends into a working area. The dotted lines are the contours for particles with amplitudes satisfying  $x^2/\sigma_x^2 + y^2/\sigma_y^2 = n^2$ . We assume  $\xi_x = \xi_y = 0.05$ . Case (a) is when the aspect ratio is  $a = 1$ , i.e. a round beam. Case (b) is when  $a = 0.1$ , i.e. a flat beam. Case (c) gives the result in the limit  $a = 0$ . (d) shows fitting the working area (shaded region) into a resonance free region in the tune space.



Ref.: W.T.Weng, *Space charge effects, tune shifts and resonances*, SLAC-PUB-4058, Aug.1986

# ( Traditional ) optimization of luminosity

◆ collision head-on or very small crossing angle

◆ space parameters  $\zeta_{x,y}$  treated as constant  
(i.e. set to saturation limit)

■  $\zeta_{x,y}$  given by bunch charge and geometry

$$\zeta_u \propto \frac{N\beta_u^*}{\sigma_u^* (\sigma_x^* + \sigma_y^*)}; \quad u = x, y$$

→ insert in luminosity formula:

Luminosity (flat beams)  $\mathcal{L} \sim E^2 \frac{\varepsilon_x}{b\beta_y^*} \zeta_x \zeta_y S$

→ large (!) emittance

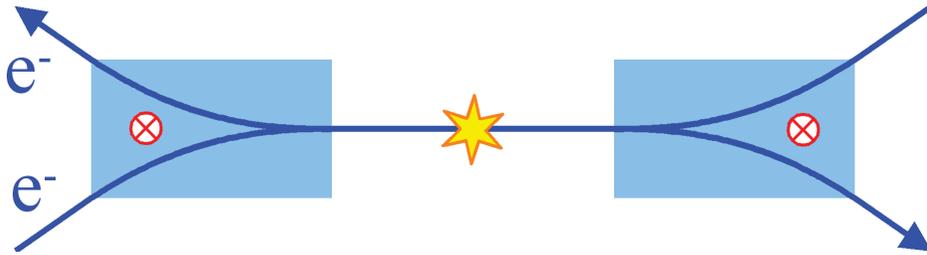
→ sharp focus – *but*

- avoid luminosity suppression due to hourglass effect:  $\beta^* > \sigma_s$  for  $S \approx 1$
- avoid longitudinal beam-beam effect: energy modulations for large  $x', y'$
- avoid tune shift increase  $\Delta Q > \zeta_{x,y}$  for  $|\Delta s| > 0$

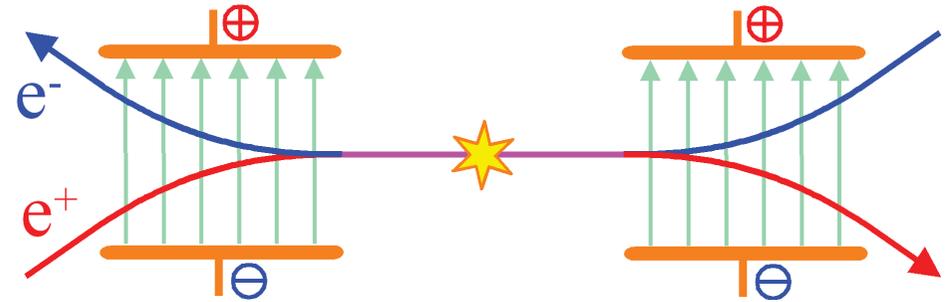
→ short distance between successive bunches

- limit:  $b = \lambda_{\text{RF}}$ , bunch distance = RF wavelength
- avoid parasitic crossings

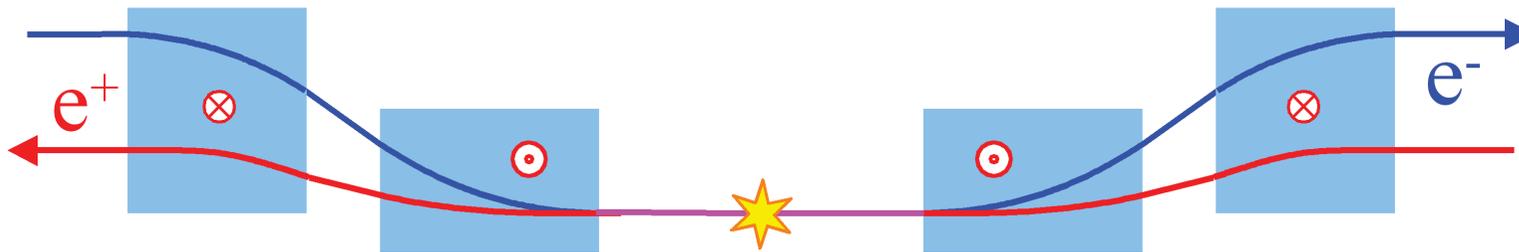
# Beam separation schemes



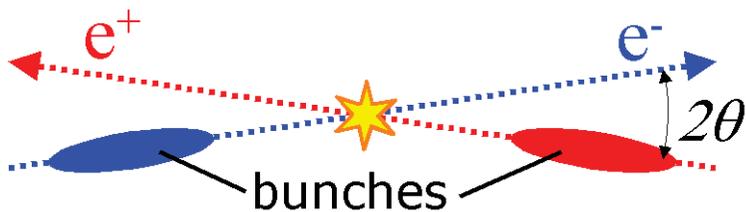
Magnetic separation for head-on collision of identical particles at same energy



Electrostatic separation for head-on collision of particle/antiparticle at same energy



Asymmetric magnetic separation for head-on collision of particle/antiparticle at different energies to boost the center of mass system.



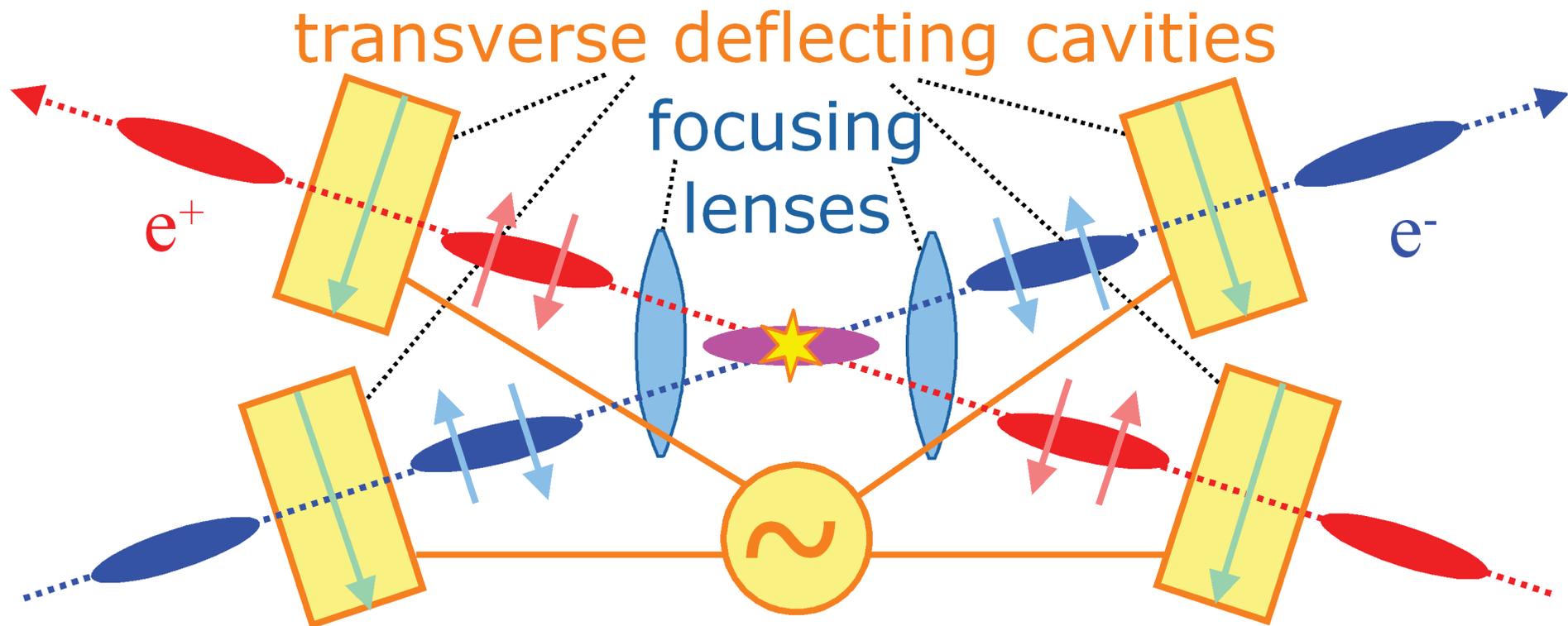
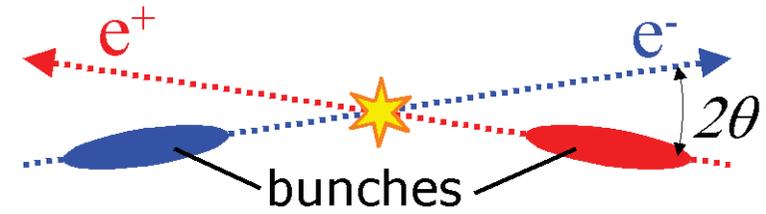
Crossing angle

- high collision frequency
- reduced luminosity due to incomplete overlap

# Crab Crossing

Crossing at angle  $\theta \ll 1$

- 😊 high bunch frequency
  - ☹️ reduced overlap/luminosity
  - ☹️ excitation of transverse/longitudinal coupling resonances
- ⇒ restore head-on collision:



In use at KEK B-factory, reached  $\mathcal{L} > 2 \cdot 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$  😊

# Crossing collision

for flat beams (  $\sigma_y \ll \sigma_x \ll \sigma_z$  )  
large crossing angle (some degree)

## 1. Reduction of interaction zone

$$2\sigma_s \rightarrow L_{\text{cross}} = 2\sigma_x / \sin \theta$$

⇒ reduction of  $\mathcal{L}$  and  $\zeta$

## 2. Low emittance beams

⇒ increase  $\zeta$  to space charge limit again

⇒  $\mathcal{L}$  increases too ⇒  $\mathcal{L}$  independent of crossing angle  $\theta$  !

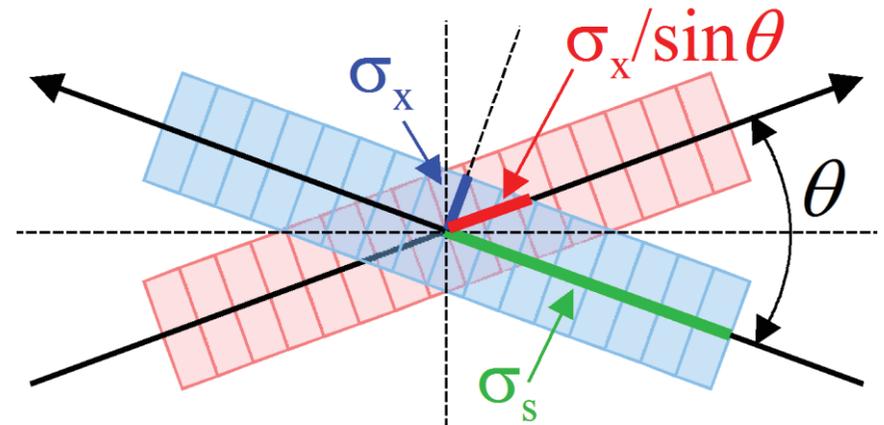
## 3. Micro-beta focusing

$$\text{Criterion: } \beta^* > \sigma_s \rightarrow \beta^* > L_{\text{cross}} \ll \sigma_s$$

⇒ sharper focusing, very small  $\beta^*$

⇒ increase of luminosity:  $\mathcal{L} \sim 1 / \beta^*$

⇒ decrease of space charge parameter:  $\zeta \sim \beta^*$



D. N. Shatilov & M. Zobov,  
ICFA BD NL 37 (2005) 99-109

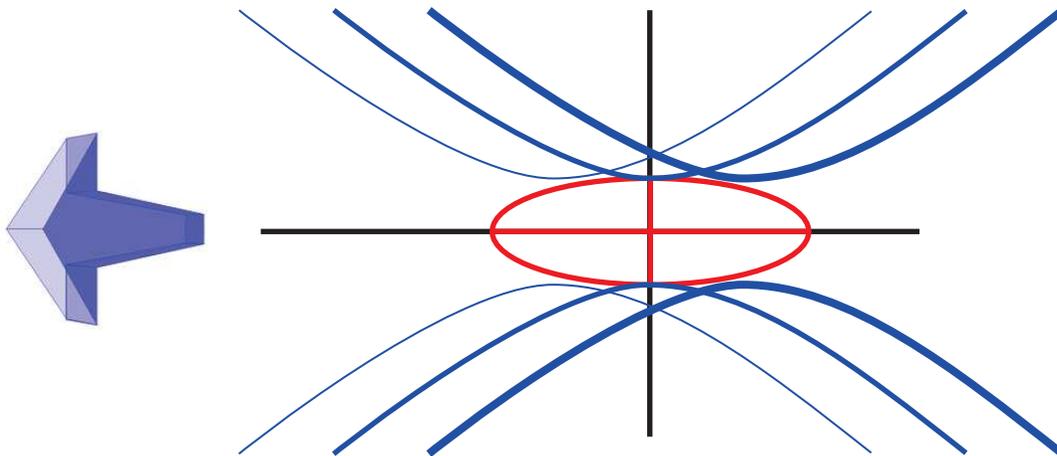
P. Raimondi, D. N. Shatilov & M. Zobov,  
EPAC-08, Genoa, 2008, WEPP045

## 4. Crab waist

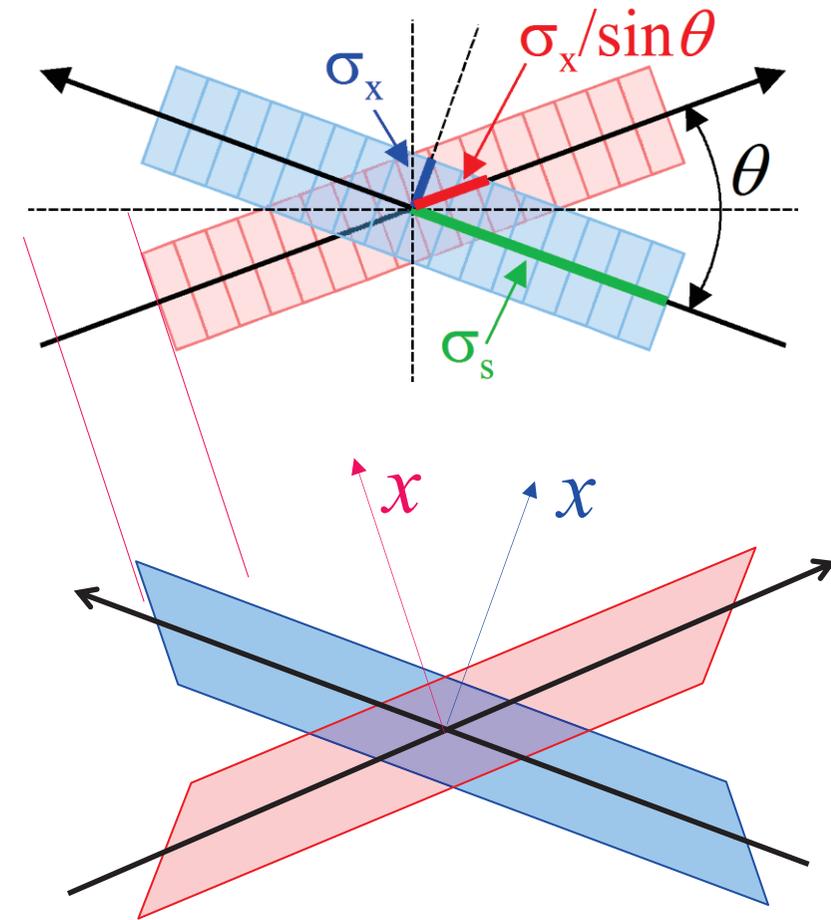
Lateral shift of focus during collisions

⇒ transverse-longitudinal coupling forces

⇒ beam instability!



colliding “blue beam” as seen from “red beam”



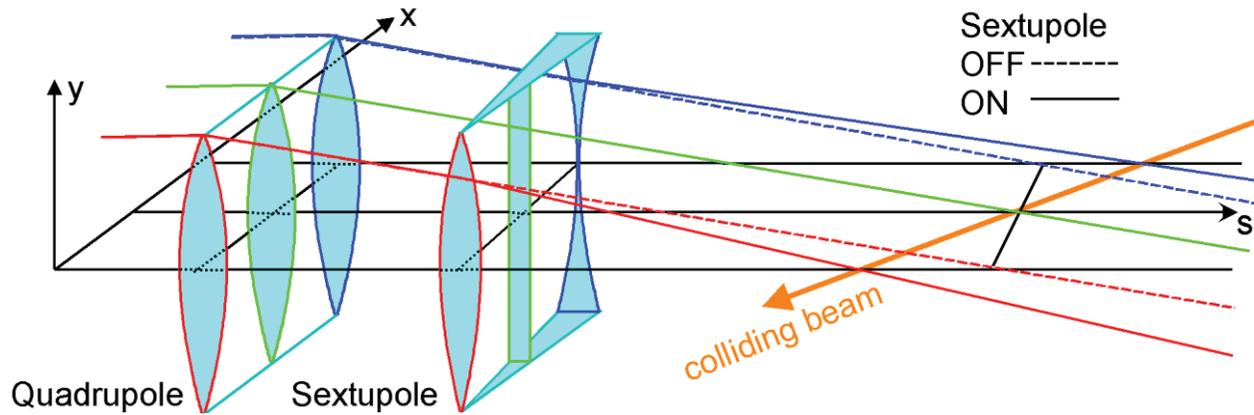
To do: adjust  $x$ -dependent  $s$ -position of focus,  
such that it coincides with the other beam's  $x$ -axis

⇒  $x$ - $s$  shear of bunch

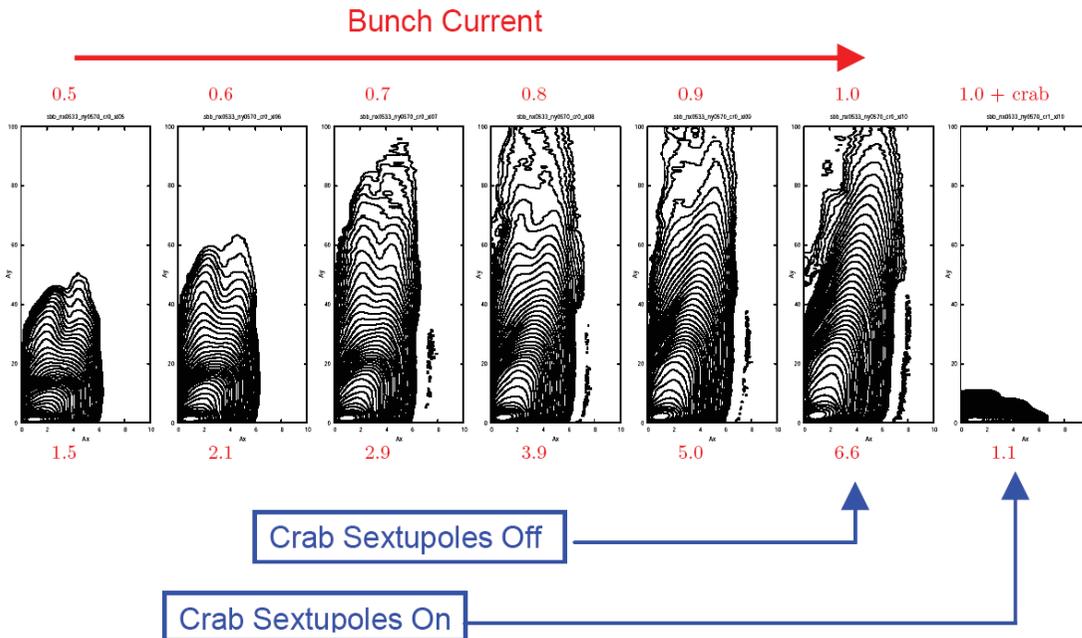


*How to do?*

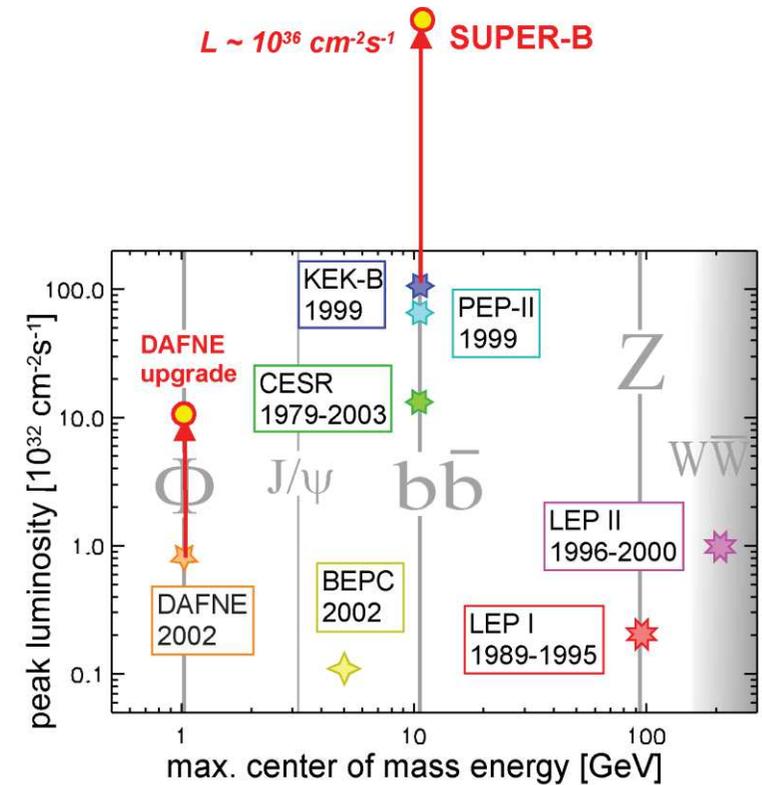
# “Crab sextupole”



## Horizontal variation of focal length



Suppression of vertical blow-up (beam instability)



Expected luminosity increase

Crab-crossing successfully tested at **DAΦNE**

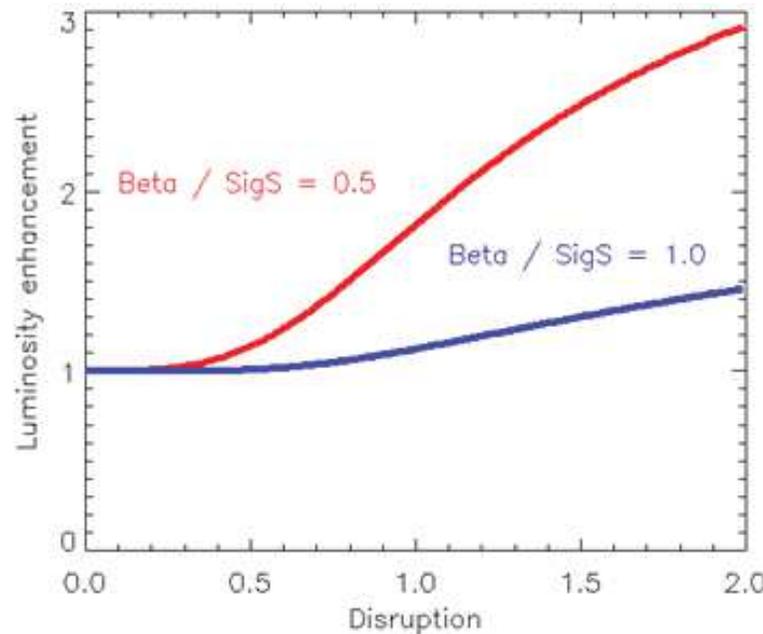
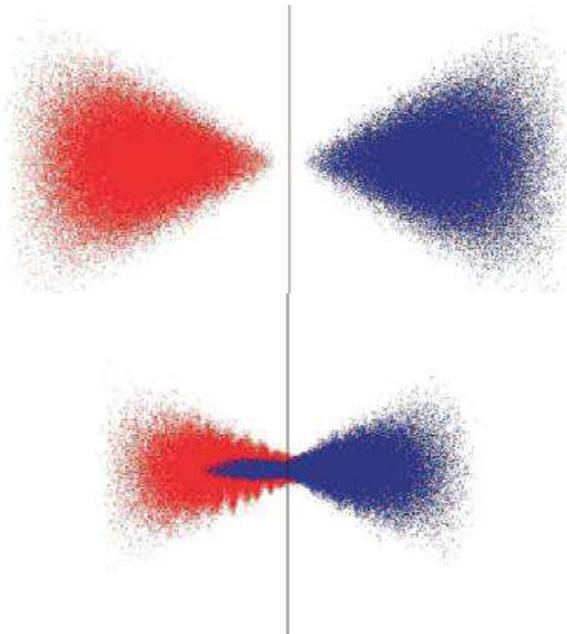
# Beam disruption

Linear collider: no re-use of beam  $\Rightarrow$  no limit on  $\zeta$ -parameter

$\Rightarrow$  beam-beam lens: focal length  $<$  bunch length,  $f < \sigma_s$

$\Rightarrow$  self-focussing (pinch-effect)  $\Rightarrow$  luminosity enhancement

$\Rightarrow$  beam disruption, parameter  $d_u = \frac{\sigma_s}{f} = \frac{2r_e}{\gamma} \frac{N\sigma_s}{\sigma_u(\sigma_x + \sigma_y)} = -4\pi\zeta_u \frac{\sigma_s}{\beta_u} \quad u = x; y$



■ M. Boronina et al., PAC-07, Albuquerque 2007, THPAN060

Luminosity enhancement

## Linear collider luminosity limitations

- sub- $\mu\text{m}$  alignment of final focus
  - **beamstrahlung** = synchrotron radiation in field of oncoming bunch
- $\Rightarrow$  photon recoil  $\Rightarrow$  undefined center of mass energy
- $\Rightarrow$  direct gamma-background
- $\Rightarrow$  pair production background